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APPENDIX G DRL AF TM NO. 58

THE DISPENSING AND BEHAVIOR OF CHAFF IN SPACE

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THE DISPENSING AND MEHAVIOR OF CHAFF IN SPACE

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J. H. Henson and J. W. Craig

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THE SYSPENSING AND BEHAVIOR OF CHAFF IN SPACE

by

J. H. Henson and J. W. Craig

I. INTRODUCTION

Chaff has been used extensively during the past two decades for establishing airborne radio scattering targets. The principal advantage of this material is the very high echo area per unit weight. Heretofore chaff employment has been confined largely to the lower portion of the earth's atmosphere. With the advent of man-made satellites and other types of space vehicles, it becomes worthwhile to consider some of the problems associated with the use of chaff above the earth's atmosphere.

II. SCOPE

This memorandum deals primarily with investigations conducted at DRL in two areas of the space-chaff problem. First, an experimental study was made to determine some possible methods of dispensing chaff at very high altitudes. Secondly, the behavior of chaff when dispensed from an earth satellite in a circular orbit was investigated.

Many uses for chaff in space have been suggested by groups throughout the country. Some possible applications include scatter communications, countermeasures, and decoys. It is not the intention of this paper, however, to deal with the uses of chaff in space or the quantities accessary for the various applications.

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III. REQUIREMENTS FOR A VERY-HIGH-ALTITUDE DISPENSER

The basic requirements for $\hat{\kappa}$ very-high-altitude dispenser are as follows:

- (1) The dispenser must separate the chaff into individual dipoles without mutilation.
- (2) The dispenser must give the dipoles pre-determined velocity with considerable precision.

For chaff to be most effective as an electro-magnetic reflector, it is necessary that the cloud be composed of individual dipoles, not groups or clusters. When chaff is dispensed from a fast-moving vehicle within the earth's atmosphere, drag forces tend to separate the chaff clusters into individual dipoles. Outside of the earth's atmosphere, the drag forces available for dipole separation are very small. Hence it is necessary to use some other means to assure efficient separation of the chaff bundle.

The velocity given to the dipoles when chaff is dispensed within the earth's atmosphere can be varied over a wide range and still produce a cloud of reasonable dimensions and dipole density. If the velocity is small, wind currents will disperse the chaff. If the initial velocity is large, the excess kinetic energy will be rapidly absorbed by the air. In any event, the cloud will retain finite dimensions for some appreciable length of time. On the other hand, chaff dispensed above the earth's atmosphere (but not in orbit) will retain the dispensing velocity indefinitely. If this velocity is large the cloud will grow to tremendous proportions and become quite diffuse in a short time. If the dispensing

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relocity is small, the bloom time" will be excessive. For chaff dispensed from a vehicle in orbit, the magnitude and direction of the dispensing velocity will determine the rate of growth of the earth chaff by t as well as the maximum dimension (width, depth, etc.) of the belt. More will be said about the behavior of orbiting chaff in later sections of this report.

*The time required for the cloud to reach useful size.

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IV. EXPERIMENTAL INVESTIGATION OF DISPENSING METHODS

A. Vapor-Pressure Method

1. General description:

This technique utilizes the rapid vaporisation and expansion of fluids to achieve the initial dipole separation and small velocity differential desired. The actual mechanics of the operation consists of soaking the dipoles thoroughly with the desired fluid and packing them into a dispensing chamber. After all excess fluid is drained off, the package is then immediately seeled under atmospheric conditions. When it is to be dispensed, the fluid-soaked chaff is thrust from the dispensing chamber into a vacuum. This vacuum causes rapid vaporisation and consequent expansion of the fluid. The vapor flow imparts velocity to the dipoles and the presence of fluid between individual dipoles ensures good dipole separation.

2. Scope of experimental work:

The experimental work was limited to observation of dipole separation and measurement of cloud velocity as a function of fluid vapor pressure and chaff type. The fluids used ranged from Dow Corning 200 milicone fluid, 1.0 centistoke viscosity with a vapor pressure at ambient temperature of 2.7 mm Hg, to acetone, which has an ambient temperature vapor pressure of 229 mm Hg.

Four types of aluminum and one of glass chaif were used in the experimental work. Aluminum chaff size ranged from $0.75 \times 0.016 \times 0.0005$ in. to $4.5 \times 0.25 \times 0.0005$ in. The length and diameter of the glass dipoles were 0.75 in. and 0.002 in. respectively.

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3. Description of test equipment:

The bell jar (volume approximately 3 ft³) and one of the ejection devices used in these investigations are shown in Figure 1. Figure 2 is a cross section drawing of the ejection apparatus and shows the method of executing chaff ejection. The chaff is exposed to the vacuum by a downward distribution of the enclosing cup. This movement is accomplished by pulling on the entry of the enclosing cup.

It should be noted that two different types of chaff containers were employed. These containers are shown in Figure 3. The upper container was used with the 3/4" length dipoles which were packed axially. The lower container was used with the +1/2" length dipoles which were packed circumferentially. Both containers were designed for radial chaff ejection.

A pictorial recording of each test was made with a Fastax camera.

4. Test procedure:

The chaff to be used in the test was bathed and cleaned thoroughly in the fluid involved. This was done to remove the chaff lacquer coating which is soluble in most of the fluids under consideration. (This procedure was not necessary for the $4-1/2^r$ chaff which has no lacquer coating)

After the bathing operation the chaff was to ked into the appropriate chaff container. It was then scaked thoroughly with fluid, the chaess fluid powed off, and the container scaled inmediately. The loaded container was then placed in the bell jar.

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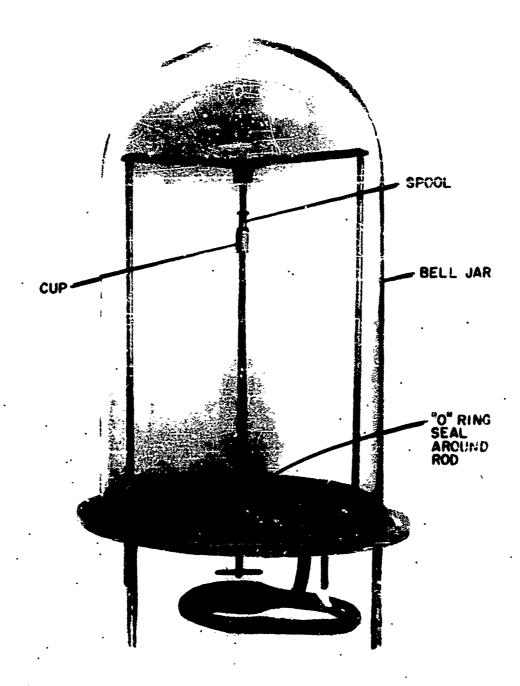


FIGURE 1
TEST SET-UP BEING USED IN VERY-HIGH-ALTITUDE
CHAFF-DISPENSING STUDIES

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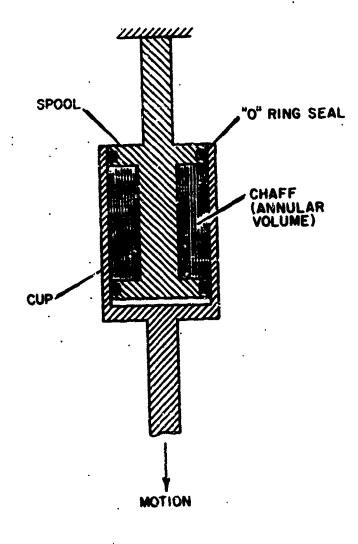
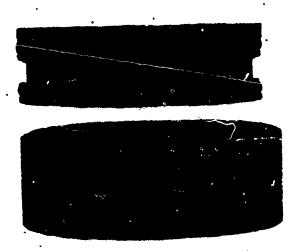


FIGURE 2
CHAFF CONTAINER FOR SIMULATED HIGH
ALTITUDE DISPENSING EXPERIMENTS

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(a) CONTAINER FOR AXIALLY PACKED CHAFF.



(b) CONTAINER FOR CIRCUMFERENTIALLY PACKED CHAFF.

FIGURE 3

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Care was taken to prevent radiant heating of the fluid and chaff from the high intensity Fastax lights during the tests. All the tests were run at bell-jar pressures of less than 0.10 mm Hg. Precautions were taken to prevent pressure leakage around the 0-ring seals of the chaff container.

5. Data reduction:

The Fastax film of the tests was first studied qualitatively with a 16 mm movie projector to observe dipole dispersion and separation. It was then studied quant latively frame-by-frame to determine dipole velocity. Basically, the velocity was determined by measuring the change in diameter of the cloud for a given number of frames. Knowing frame speed and an appropriate dimensional scale factor, the average dipole velocity could be readily calculated.

. 6. Results:

The principal results of the experimental investigations are given in Figure 4. Note that the dipole velocity is approximately linear with $\frac{PA}{R}$. In plotting these data, several very irregular points have been omitted.

 $\mbox{\bf A}$ complete description of all tests is given in tabular form in Appendix A.

The results of Figure 4 stated in equation form are:

$$V = 2.03 + 0.0394 \sqrt{\frac{pA}{n}}$$

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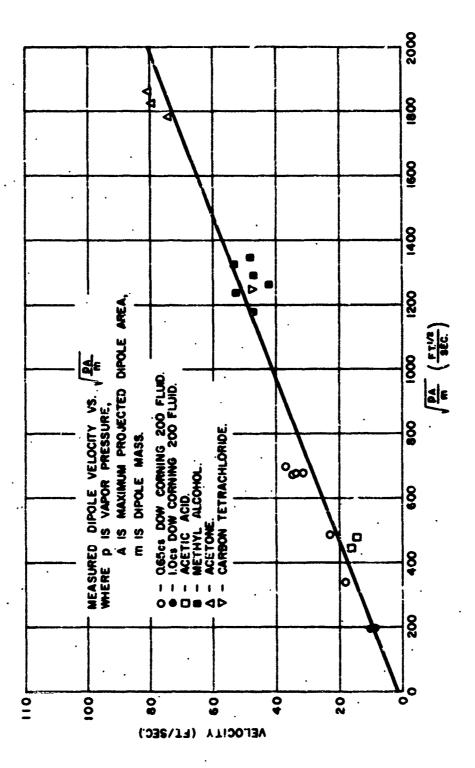


FIGURE 4
EXPERIMENTAL RESULTS

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The coefficients were determined from a least squares fit of the experimental points in Figure 4. It should be stressed that this is strictly an empirical statement. It is approximate, and is known to hold only for the vapor-pressure range of the table in Appendix A. It should be used with caution outside this range.

Figure 5 is a print of several frames of film from Test K which employed methyl alcohol as the ejecting fluid. Figure 6 (2 pages) is a similar print of Test W which employed no fluid. Note the great difference in the behavior of the two chaff clouds.

The frame speed for both runs was approximately 2300 frames/sec, but only every tenth frame was printed. The horizontal bar appearing in Figures 5 and 6 was used for calibration in the determination of dipole velocity.

7. Limitations:

The use of fluid vaporisation to dispense chaff is limited to very low pressure or high-altitude conditions. It is necessary that the ambient pressure be very low compared to the vapor pressure of the fluid in use.

Since vapor pressure varies radically with fluid temperature, knowledge of temperature conditions at the time of dispensing is necessary. The unevaporated fluid is cooled during dispensing as vaporization takes place. However, the conditions in the test chamber and in space are essentially the same in this respect. The radiation loss to space during the very short dispensing period is negligible.

Dipole separation may be incomplete with a fluid of relatively low vapor pressures. In the tests using 1.0 centistoke viscosity Dow

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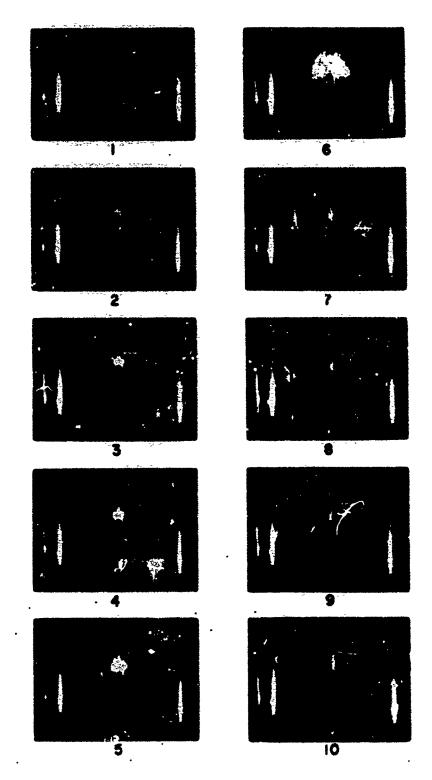


FIGURE 5
CHAFF EJECTION USING METHYL ALCOHOL AS EJECTING FLUID

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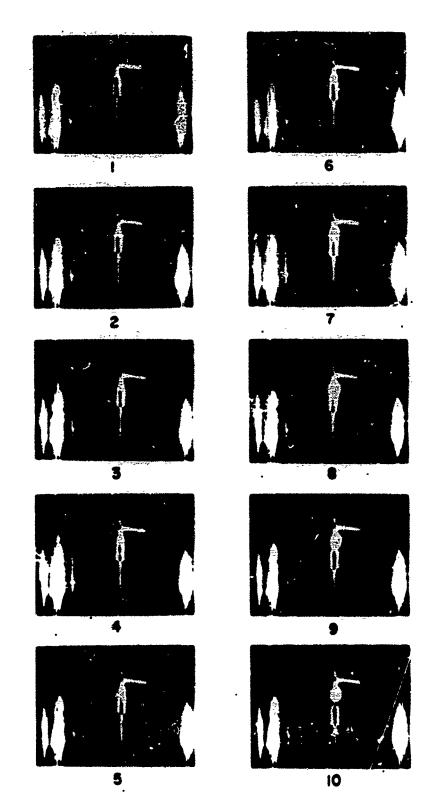


FIGURE 6

CHAFF EJECTION USING NO EJECTING FLUID

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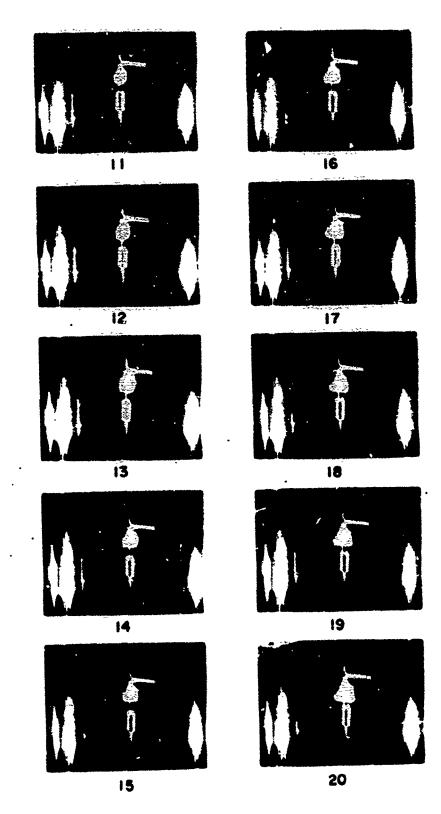


FIGURE 6 (CONTINUED)

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Corning 200 Silicone fluid (wapor pressure about 2.7 mm Hg at 72°F), some dipoles stuck together in clumps of three or four.

This method of dispensing is not applicable to low-frequency "rope" type chaff. Other specialized methods of dispensing will be necessary for this variety of chaff.

B. Low-Pressure Cas Dispenser

Another method of chaff dispensing investigated involved a low-pressure gas dispenser. This dispenser consists of a glass cylinder packed with chaff and pressurized with low-pressure air. The chaff is dispensed by shattering the glass with bullets fired by high-pressure air as shown in Figure 7.

Figure 8 shows one of the dispensers designed and built at DRL. Several units were test 3 but sealing difficulties, particularly at the ends, were encountered. The project was discontinued in favor of concentration on the vapor-pressure method.

C. Spin Dispensers

The spin dispenser is a mechanical device for ejecting and dispensing the chaff, as shown in Figure 9. Three dispensers of this nature were designed and built at Defense Research Laboratory and installed as "piggyback" equipment in a Thor missile nose-cone. The missile was fired from Cape Canaveral in January of 1960. For a complete description of this experiment and the associated equipment see DRL AF Technical Memorandum No. 50 (Contract AF 33(616)-5164).

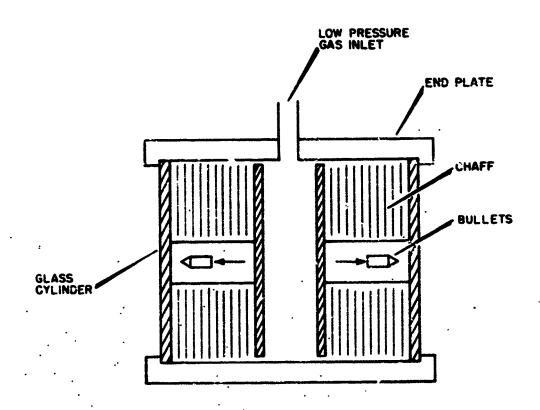


FIG. 7
LOW-PRESSURE GAS CHAFF EJECTION SETUP

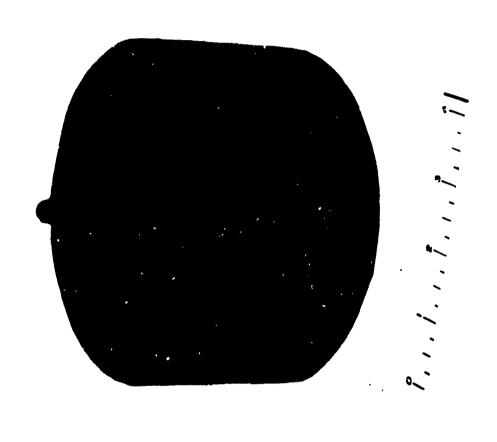
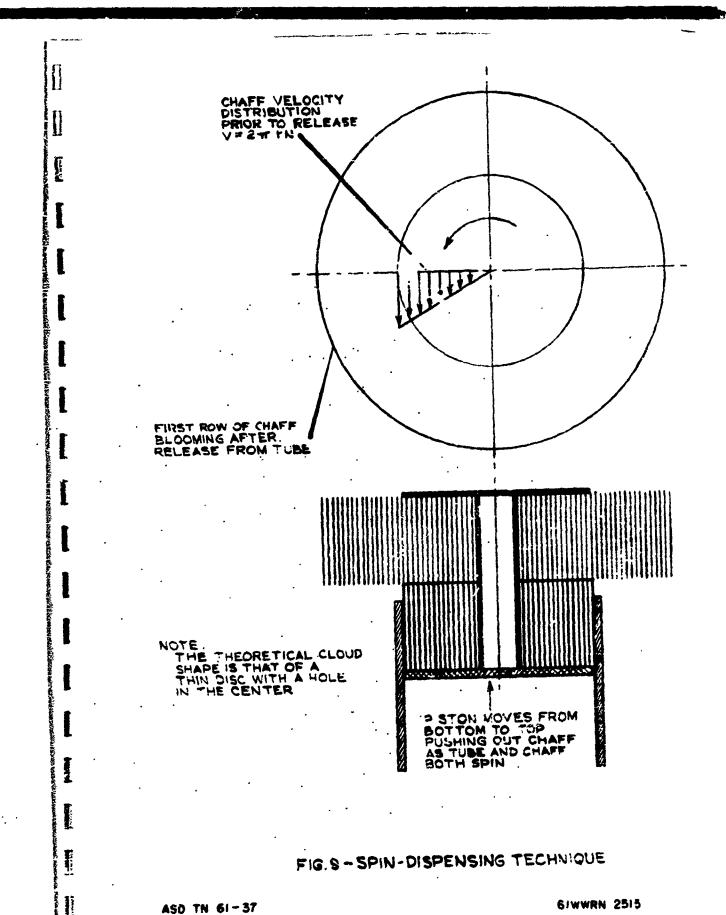


FIGURE 8 CHAFF CONTAINER FOR LOW PRESSURE GAS EJECTION

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V. THEORETICAL INVESTIGATIONS OF CHAFF BEHAVIOR IN SPACE

A. Scope of Study

Theoretical studies have been limited to the behavior of chaff when dispensed from a vehicle in a circular geocentric orbit. From curves and other information included herein, it is easy to determine the effect of the dispensing velocity on certain pertinent chaff cloud parameters. Of particular interest is the time required to form a complete belt of chaff around the earth.

B. Assumptions

In the theoretical analysis the following assumptions were made:

- (1) Drag on the dipoles is negligible during the period of interest. This will be true if the orbit is sufficiently high above the earth.
- (2) The dispensing velocity is imported to the dipoles instantaneously. Experimental results from the vapor-pressure study show that the velocity is obtained within a few milliseconds.
- (3) The simplified two-body equations are sufficiently accurate to describe the orbit of the dipoles. This will be true so long as the orbit is near enough to the earth that the attraction to the moon, sun, and planets can be neglected. The effects of earth oblateness are also neglected.
- (4) Photoelectric and magnetic effects are negligible.

C. The Elliptical Orbit

Figure 10 shows an ellipse and defines some of its important properties. The path of a particle in orbit about the earth is an ellipse with the earth at one focus. A circular orbit is a special ellipse for which the eccentricity is zero. (A discussion of orbits can be found in any good book on mechanics.)

Some of the more important relationships for simple two-body systems are given below.

(1)
$$a = \frac{4}{2E}$$

where
$$\mu = g_0 R^2 = constant$$

(2)
$$e = \sqrt{1 + \frac{2Bh^2}{\mu^2}}$$

g_o= Acceleration of gravity at surface
 of earth-lt/sec

(3) E = 1/2 V² - $\frac{\mu}{r}$ (Constant for a particular orbit)

R = Redius of earth - ft

(4) $V^2 = \mu \left(\frac{2}{r} - \frac{1}{a} \right)$

E = Sum of kinetic and potential energy per unit wass (E<O for ellipse)

h = Angular momentum per unit mass

(5) $\tau = \frac{2\pi\mu}{\sqrt{(-2E)^3}}$

τ = periodic time - sec

D. Dispensing Effects

1. Preliminary Consideration

Consider a vehicle in a circular geocentric orbit capable of dispensing chaff with a velocity ΔV relative to itself in all directions simultaneously. (See Figure 11.)

Throughout this discussion, the subscript 0 wherever used refers to the conditions of the original circular orbit. The subscript 1 refers to the orbital conditions produced by ΔV_{η} and so forth.

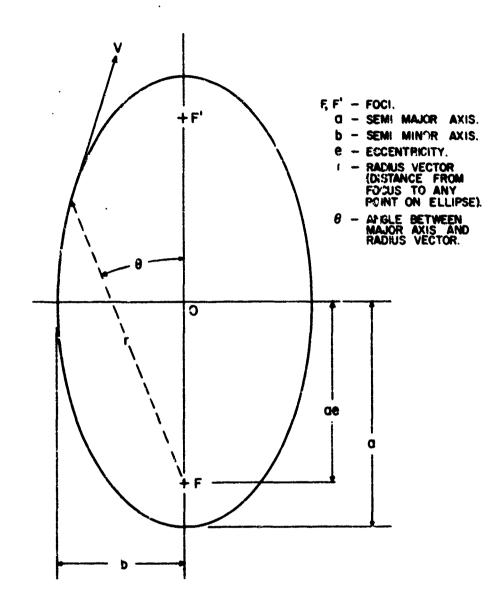
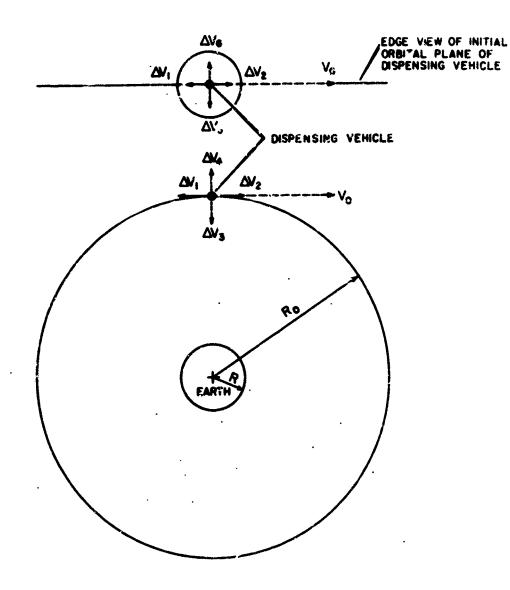


FIGURE 10
THE ELLIPTICAL ORBIT

£ 1_



Vo = INITIAL CIRCULAR ORBITAL VELOCITY.

Ro = MITIAL ORBITAL RADIUS.

R = RADIUS OF EARTH = 20.925 X 10^6 FT. ΔV = PERTURBATION OR DISPENSING VELOCITY. $|\Delta V_1| = |\Delta V_2| = |\Delta V_3| = |\Delta V_4| = |\Delta V_5| = |\Delta V_6|$

FIGURE II LIMITING DISPENSING CONDITIONS

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its total velocity vector changed in magnitude or direction or both. Each dirole will theoretically be transferred to a unique elliptical orbit.

It is difficult to define as a function of time the shape and size of a group of orbiting dipoles when dispensed in the manner described in the above paragraph. In order to simplify the problem, only the dimensional limits of the chaff cloud were studied. The AV's of the directions shown in Figure 11 establish the maximum dimensions (width, depth, etc.) of the chaff cloud. It is now worthwhile to consider in datail the effect of these limiting AV's on the orbit parameters.

2. AV and AV 2

The time required to "belt" the earth as well as the maximum depth of the belt are determined by δW_1 and δW_2 . Figure 12 shows qualitatively the effect of δW_1 and δW_2 with respect to the initial circular orbit. A dipole ejected with velocity δW_1 will have a total velocity (immediately following ejection) and energy per unit mass less than that of the dispensing vehicle. On the other hand, δW_2 increases the dipole velocity and energy per unit mass. The vector addition of V_0 to the various δW_2 's shows clearly that δW_1 and δW_2 produce the extremes in dipole energy change. Since the periodic time is a function of E only, δW_1 and δW_2 also produce the extremes in periodic time (See Equation 5). It is these extremes in the period that determine the time required for the chaff cloud to "belt" the earth. It should be noted that the orbital plane of dipoles dispensed with δW_1 and δW_2 is the same so that of the dispensing vehicle.

Using Figures 13 and 14, it is possible to determine the number of orbits as well as the time required to form a complete earth belt

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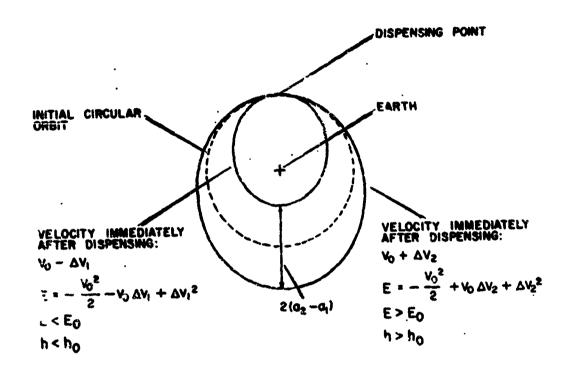
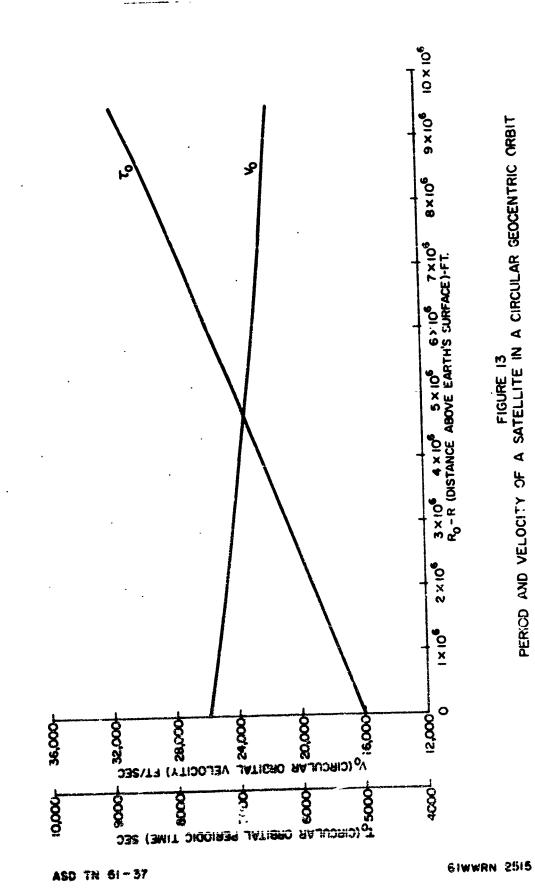
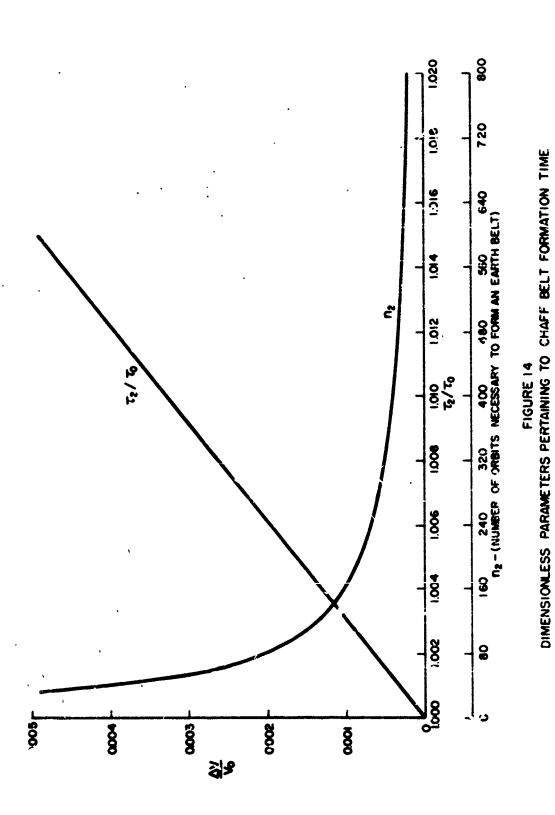


FIGURE 12

QUALITATIVE ORBITS AFTER PERTURBATION
FOR AND AV2



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when the chaff is dispensed with ΔV limits shown in Figure 11. It is interesting to note that n (Figure 14) is a function only of $\frac{\Delta V}{V}$ and not R_0 .*

The use of these curves can best be explained by an example.

Given

AW = 30 ft/sec (determined by fluid vapor pressure)

$$R_0 - R = 6 \times 10^6$$
 feet

From Figure 13 for circular orbits

Then
$$\frac{2V}{V_0} = \frac{30}{22,800} = 0.00132$$

From Figure 14

Time required to belt the earth

$$= n_2 \tau_2 = \frac{7440 (125)}{3600} = 258 \text{ hours}$$

From Equation 1 it is readily seen that M_1 and M_2 also determine the maximum depth of the chaff belt. This thickness is given

*For a proof of this and for various derivations, see Appendix B.

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by 2 ($a_2 - a_1$) as shown in Figure 12: This parameter is plotted in Figure 15 as a function of $\frac{M^{-4}}{V}$. Again the use of this curve can best be explained by an example.

Given (from above example)
$$\frac{\Delta V}{V_o} = 0.00132$$

$$R_o = 6 \times 10^6 + 20.925 \times 10^6 \text{ ft}$$
From Figure 15
$$\frac{2(a_2 - a_1)}{R_o} = 0.0105$$

$$2(a_2 - a_1) = 0.0105 (26.925 \times 10^6)$$

$$= 28.25 \times 10^4 \text{ ft}$$

$$= 53.5 \text{ miles}$$

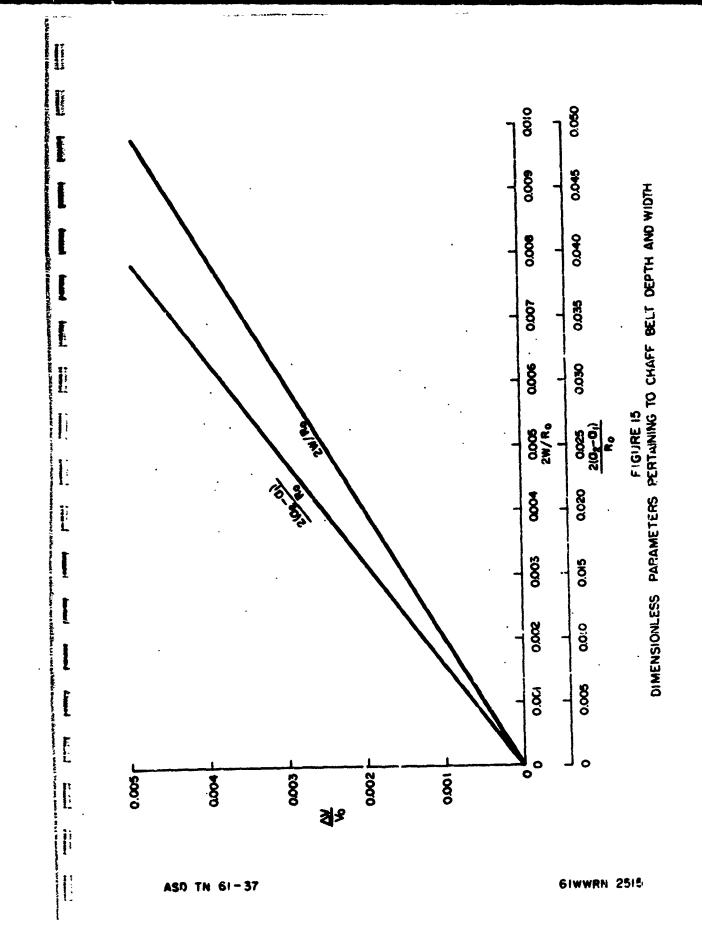
3. My and My

Figure 16 shows qualitatively the result of N_3 and N_4 with respect to the initial circular orbit. Note that the apse lines of these two orbits are shifted by plus and minus 90° from the dispensing point.** It is easily shown from energy considerations that these orbits do not produce the maximum cloud depth nor the extremes in orbital period. The orbital plane of dipoles dispensed with N_3 and N_4 is the same as that of the dispensing vehicle.

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^{*} See Appendix B for derivations. **See Appendix B for proof.



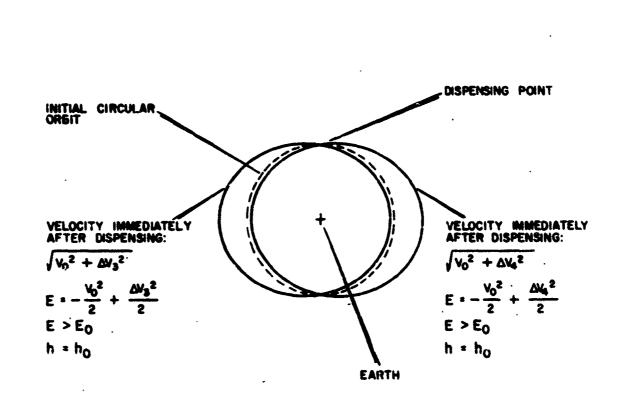


FIGURE 16 QUALITATIVE ORBITS AFTER PERTURBATION FOR ΔV_3 AND ΔV_4

4. ΔV_5 and ΔV_6

The maximum width of the chaff belt is determined by ΔV_5 and ΔV_6 . Figure 17 shows qualitatively the effect of these dispensing velocities. The principal result is to rotate the planes of the orbits, although other orbital characteristics are also changed. However, from simple energy considerations it is readily seen that ΔV_5 and ΔV_6 do not produce the maximum depth of the belt nor the extremes in orbital period.

From the lower curve of Figure 15 it is possible to determine the maximum width of the belt.* An example will readily explain the use of this curve.

Given (from previous examples)

From Figure 15

$$\frac{2N}{R_0} = .0026$$

- 70,000 ft

* 13.2 miles

*See Appendix B for derivation of the equation.

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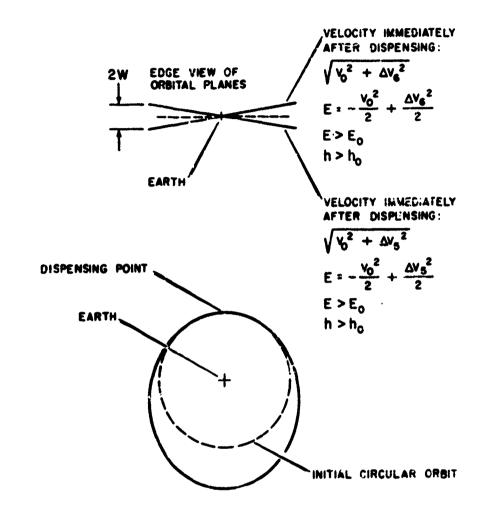


FIGURE 17 QUALITATIVE ORBITS AFTER PERTURBATION FOR ΔV_{S} AND ΔV_{S}

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5. Other Considerations

It is interesting to consider the result obtained when the chaff is dispensed from the orbital vehicle without a velocity component parallel to the original velocity vector. (All ΔV 's are equal and perpendicular to the original velocity vector.) After dispensing, all dipoles will have the same energy and hence the same period (Equation 5). Instead of forming a belt, the chaff cloud will grow and change as a function of time in some peculiar cyclic manner, repeating this cycle each orbit. All the chaff will theoretically arrive back at the dispensing point at the same time. Thus as ΔV_1 , ΔV_2 , and all other velocity components parallel to V_0 approach zero, the number of orbits required to produce an earth helt approaches infinity, as indicated by Figure 14.

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VI. CONCLUSIONS

- (1) The vapor-pressure technique is very effective for separating dipoles in a space-like environment.
- (2) When dispensing fluid-saturated chaff in a low-pressure environment, the dipole velocity is approximately linear with the square root of the fluid vapor pressure.
- (3) Chaff dispensed candidrectionally from a vehicle in a circular geocentric orbit will form a belt around the earth.
- (4) Chaff given a uniform dispensing velocity perpendicular to the original circular orbital velocity will produce a chaff cloud which grows and changes in some cyclic manner. The cycle will repeat once each orbit and the cloud theoretically will not form an earth belt.

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APPENDIX B

Definition of Terms and Nomenclature used in Text and Appendix A NOTE: Subscripts 1 through 6 refer to ΔV 's 1 through 6 as defined by Figure 11 of text.

- μ Generativic orbital constant = $g_0 R^2 = 140.99187 \times 10^{14} \text{ ft}^3/\text{sec}^2$
- τ Orbits... period
- E Total specific energy of orbit
- V_{Ω} Dipole or vehicle velocity in the original circular orbit
- V Dipole velocity after ejection
- R Radius of cartn
- $R_{_{\rm C}}$ Radius of original circular orbit measured from the earth's center
- ΔV Velocity perturbation relative to the dispensing vehicle
- n Number of orbits necessary to form a chaft belt around the each
- a Semi-major axis of elliptical orbit
- r Radius of a point in elliptical orbit
- θ Angular orientation of a point in orbit measured counterclockwise from the apse line of the orbit
- e Eccentri ing of orbit
- h Specific angular momentum
- b Semi -minor axis of elliptical orbit
- W Maximum cloud width
- ϕ Angle becomes the orbits of dipoles dispensed with ΔV_{5} and ΔV_{6} and the original orbit

I. Derivation of n Expression

(1)
$$\tau_1 = \frac{2\pi\mu}{(-2E_1)^{\frac{5}{2}/2}}$$

and

(2)
$$\tau_2 = \frac{2\pi\mu}{(-2E_2)^{\frac{3}{2}/2}}$$

where

(3)
$$R_1 = 1/2 V_1^2 - \frac{\mu}{R_0}$$

and

(4)
$$v_1 = v_0 - \Delta v_1$$

(5)
$$E_2 = 1/2 V_2^2 - \frac{\mu}{R_0}$$

But:

(1)
$$V_0^2 = \frac{\mu}{R_0}$$

Applying (4), (6), and (7) to (3) and (5) and noting that $|\Delta V_1| = |\Delta V_2| = \Delta V$ we obtain:

(3')
$$E_1 = \frac{-v_0^2 - 2v_0\Delta V + \Delta V^2}{2}$$

(5') $v_2 = \frac{-v_0^2 + 2v_0\Delta V + \Delta V^2}{2}$

Note: from (3) and (5) that $\left| E_2 \right| < \left| E_1 \right|$

Therefore:

τ₂ > τ₁

At the instant the earth chaff belt is complete

Solving (8) for n_2 .

(9)
$$n_2 = \frac{\tau_1}{\tau_2 - \tau_1}$$

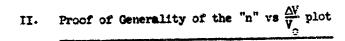
Then substituting (1) and (2) into (9) and simplifying we obtain:

(9')
$$r_2 = \frac{1}{\left(\frac{E_1}{E_2}\right)^{3/2} - 1}$$

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Consider two sets of dipoles denoted by subscriets a and b. Assume they are ejected into orbit such that:

 $\Delta V_a = k \Delta V_b$ and

 $\frac{\Delta V_a}{V_{co}} = \frac{\Delta V_b}{V_{ch}}$ or

Then five (9') of I in the appendix:

$$n_{a} = \frac{1}{\left(\frac{E_{la}}{E_{ca}}\right)^{3/2} - 1}$$

(2)
$$n_b = \frac{1}{\left(\frac{R_{1b}}{R_{2b}}\right)^{3/2} - 1}$$

The number of orbits required to belt the earth $\cdot s$ determined by ΔV_{γ} and ΔV_2 . From equations (3') and (5') of I in the appendix:

$$E_{1a} = \frac{-v_{oa}^{2} - 2 v_{oa} \Delta v_{a} + \Delta v_{a}^{2}}{2}$$

(3)
$$E_{2a} = \frac{-V_{0a}^{2} + 2 V_{0a} \Delta V_{a} + \Delta V_{a}^{2}}{2}$$

$$\mathbf{E}_{1b} = \frac{-\mathbf{V}_{ob}^2 - 2 \, \mathbf{V}_{ob} \, \Delta \mathbf{V}_b + \Delta \mathbf{V}_b^2}{2}$$

(4)
$$R^{5p} = \frac{-\Lambda^{0p}}{-\Lambda^{0p}} + 5\Lambda^{0p} \nabla \Lambda^{p} + \nabla \Lambda^{p}$$

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$$E_{1b} = k^{2} \frac{\left[-V_{0a}^{2} - 2 V_{0a} \Delta V_{a} + \Delta V_{a}^{2}\right]}{2} = k^{2} E_{1a}$$

$$(4')$$

$$E_{2b} = k^{2} \frac{\left[-V_{0a}^{2} + 2 V_{0a} \Delta V_{a} + \Delta V_{a}^{2}\right]}{2} = k^{2} E_{2a}$$

Substitute (4') into (2):

(2')
$$n_b = \frac{1}{\left(\frac{k^2 E_{1a}}{k^2 E_{2a}}\right)^{3/2} - 1} = \frac{1}{\left(\frac{E_{1a}}{E_{2a}}\right)^{3/2} - 1}$$

or $n_b \propto n_a$ for $\frac{\Delta V_a}{V_{oa}} = \frac{\Delta V_b}{V_{ob}}$

III. Derivation of Equation for Dimensionless Maximum Cloud Depth Parameter

From basic central force orbit equations:

(1)
$$a_2 = \frac{71}{2R_2}$$

(2)
$$a_1 = \frac{-\mu}{2E_1}$$

Then the maximum depth of separation between the two orbits is given by:

or
$$2(a_{2} - a_{1}) = -\mu \left(\frac{1}{E_{2}} - \frac{1}{E_{1}}\right)$$
$$2(a_{2} - a_{1}) = 2\mu \left(\frac{1}{V_{0}^{2} - 2V_{0}\Delta V - \Delta V^{2}} - \frac{1}{V_{0}^{2} + 2V_{0}\Delta V - \Delta V^{2}}\right)$$

(3)
$$5(a^5 - a^1) = \frac{(A^0 - 5A^0 - VA_5) (A^0 - VA_5)}{8^{17} A^0 - VA_5}$$

Now note that:

$$R_0 = \frac{\mu}{v_0^2}$$

Then the dimensionless maximum cloud depth parameter is defined as:

and is given by:

(4)
$$\frac{2(a_2 - a_1)}{R_0} = \frac{8}{\frac{V_0}{\Delta V} - \frac{6\Delta V}{V_0} + \left(\frac{\Delta V}{V_0}\right)^3}$$

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IV. Proof of Apse Line Shift for Dipoles Dispensed with ΔV_3 and ΔV_4

The parametric equation for an elliptical conic section is given by:

$$(1) \quad \frac{\mathbf{r}}{\mathbf{a}} = \frac{1 - \epsilon^2}{1 + \epsilon \cos \theta}$$

or solving for $\cos \theta$:

(1')
$$\cos \theta = (1 - \epsilon^2) \frac{a}{r\epsilon} - \frac{1}{\epsilon}$$

But € is given by:

$$(2) \quad \epsilon = \sqrt{1 + \frac{210h^2}{u^2}}$$

and: (3)
$$a = \frac{\mu}{2E}$$

Substitute (2) and (3) into (1').

(1")
$$\cos \theta = \frac{\frac{h^2}{\mu} - r}{r\sqrt{1 + \frac{2Eh^2}{\mu^2}}}$$

Then at the instant of velocity perturbation:

(4)
$$r = R_0 = \frac{\mu}{V_0^2}$$

and for dipoles dispensed with $\Delta V_{\frac{1}{3}}$ and $\underline{\Delta V}_{\frac{1}{4}}$:

(5)
$$\left|\Delta V_{3}\right| = \left|\Delta V_{1}\right| = \Delta V$$

(6)
$$v_3^2 = v_4^2 = v_2^2 + \Delta v^2$$

(7)
$$E_3 = 1/2 V_3^2 - \frac{\mu}{R_0}$$

(8)
$$\mathbf{E}_{4} = 1/2 \, \mathbf{V}_{14}^{2} - \frac{\mu}{\mathbf{R}_{0}}$$

Applying (4), (5), (6) to (7) and (8) we find:

(9)
$$E^2 = E^{\dagger} = \frac{5}{-\Lambda_5 + \nabla \Lambda_5}$$

Since the angular momentum is unchanged by the velocity perturbation

(10)
$$h = V_0 R_0 = \sqrt{\mu R_0}$$

Substitute (4), (9), and (10) into (1"):

$$(1''') \cos \theta = \frac{\frac{\mu R_0}{\mu} - R_0}{R_0 \sqrt{1 + 2 \frac{(-V_0^2 + \Delta V^2) V_0^2 R_0^2}{2 \mu^2}}}$$

or
$$\cos \theta = \frac{0}{\Delta V} = \frac{0}{V_0}$$

or
$$\theta = 290^{\circ}$$

V. Derivation of Equation for Dimensionless Maximum Cloud Width Parameter

Maximum cloud width may be expressed by:

where

(2)
$$\varphi = \operatorname{Ten}^{-1} \frac{\Delta V}{V_o}$$

where

$$\Delta V = |\Delta V_5| = |\Delta V_6|$$

(3) b =
$$e^{\sqrt{1-e^2}}$$

$$(4) \quad \epsilon \quad = \sqrt{1 + 2 \frac{\mathbf{E} h^2}{\mu^2}}$$

(5)
$$E = \sqrt{V_0^2 + \Delta V^2} R_0$$

(6)
$$E = 1/2 \left(v_o^2 + \Delta v^2 \right) - \frac{\mu}{R_o}$$

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$$E = \frac{2}{\sqrt{5} + \nabla \Lambda_5}$$

(7) $a = \frac{-\mu}{2E} = \frac{\mu}{V_0^2 - \Delta V^2}$

Substitute (5) and (6) into (4).

$$(4') \in = \left(\frac{\nabla}{\nabla}\right)^2$$

Substitute (7) and (4') into (3).

(3')
$$b = \frac{\mu}{V_0^2} \sqrt{\frac{V_0^2 + \Delta V^2}{V_0^2 - \Delta V^2}}$$

Also note that:

(8)
$$\sin \varphi = \frac{\Delta V}{\sqrt{V_0^2 + \Delta V^2}}$$

Substitute (3') and (8) into (1).

$$(1') 2W = \frac{2\mu\Delta V}{V_0^2 \sqrt{V_0^2 - \Delta V^2}}$$

Define the dimensionless maximum cloud width parameter ...:

Dividing (1') by R_c, we obtain:

(9)
$$\frac{2M}{R_0} = \frac{2\Delta V}{\sqrt{V_0^2 - \Delta V^2}}$$